

Applications of UCN in fundamental particle physics:

- Search for non-zero neutron electric dipole moment,
- Measurement of the neutron lifetime,
- Experiments with the gravitationally bound quantum states of neutrons,
- Search for non-zero neutron electric charge,
- ...

All these experiments are:

- Strongly limited by low UCN density (neutron electric dipole moment, neutron charge, gravitational states);**
- Or, Systematical errors would be easier revealed if higher neutron density were available (neutron lifetime, neutron electric dipole moment).**

Possible approaches to increase UCN density:

- Solid deuterium neutron converter in vicinity of a cold-neutron (optimum $\sim 5\text{\AA}$) source – inside a reactor or a spallation source;
- Liquid helium super-thermal source at an external beam of cold neutron (optimum $\sim 10\text{\AA}$);
- Equilibrium cooling of very cold neutrons (optimum $\sim 50\text{\AA}$) at ultracold nanoparticles.

G.Pignol, K.V.Protasov, V.V.N. – arXiv:nucl-th/0510021 6 Oct 2005

- 1. Model of free nanoparticles**
- 2. Model of the infinite moderator**
- 3. Realistic moderators**

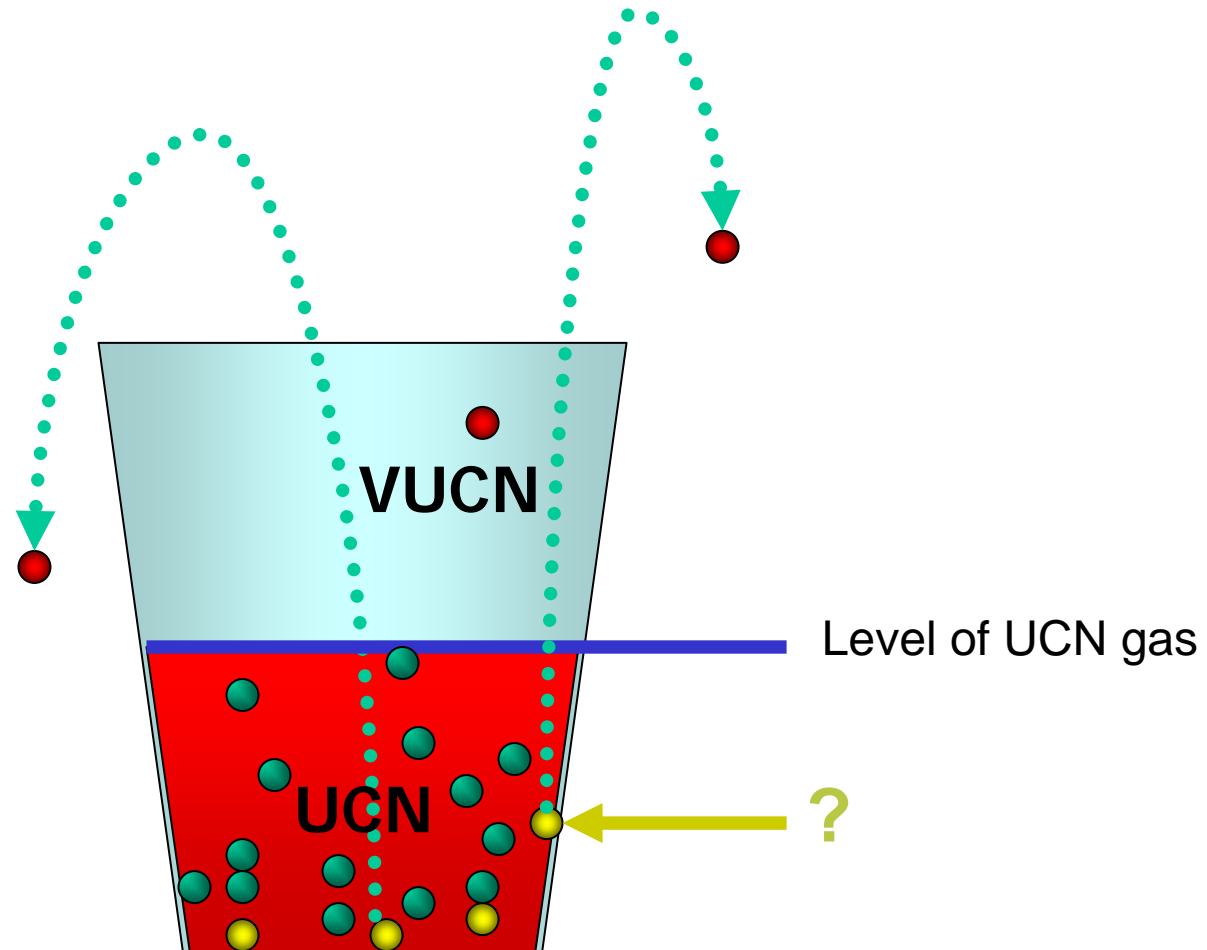
Systematics for neutron lifetime measurements ?

$VUCN$: ●

$$\Delta E \sim 10^{-7} \text{ eV}$$

$$\Delta V \sim 1 \text{ m/s}$$

$$P_{VUCN} \sim 10^{-8} - 10^{-3}$$



A.V.Strelkov et al, NIM 440A(3), 695-703 (2000)

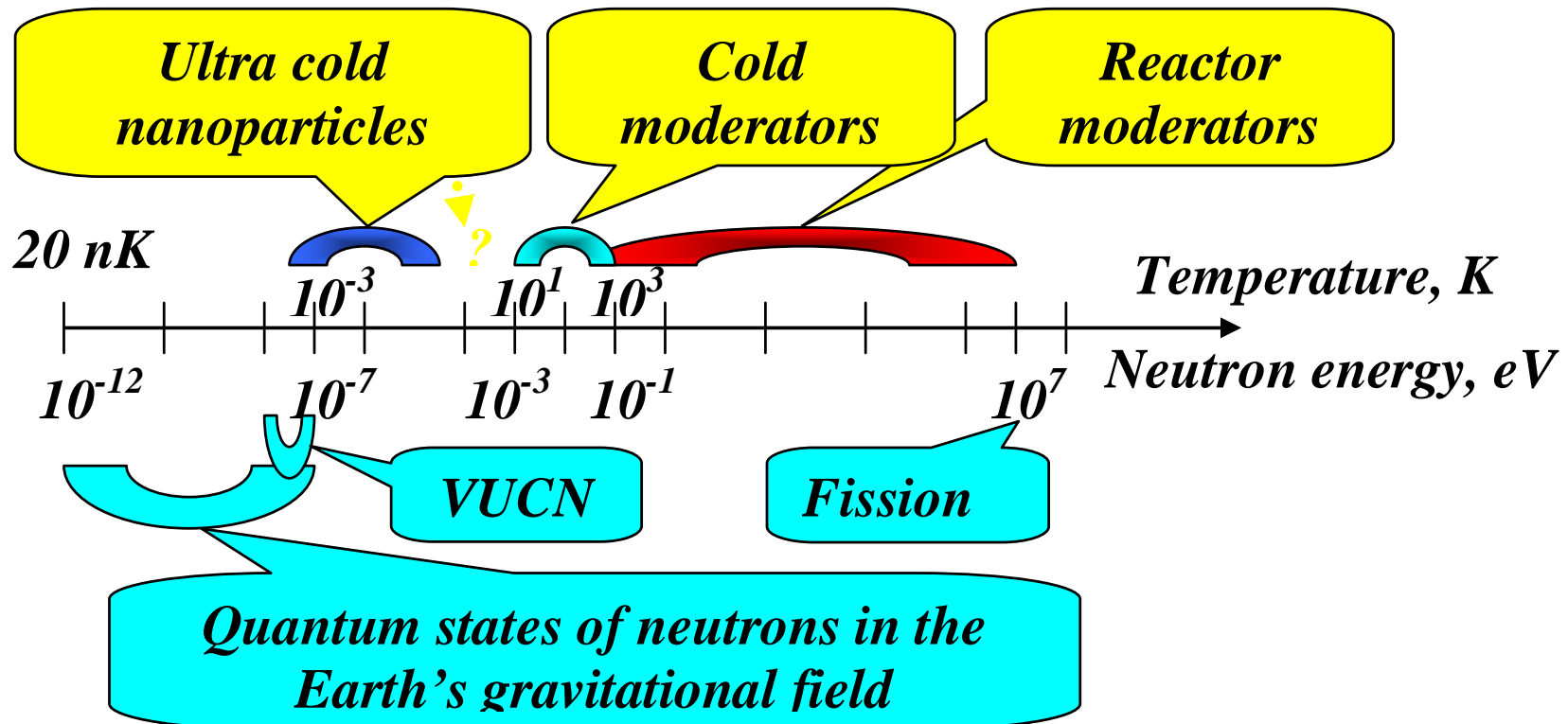
D.G.Kartashov et al, Physics of Atomic Nuclei 65(11), 1-4 (2002)

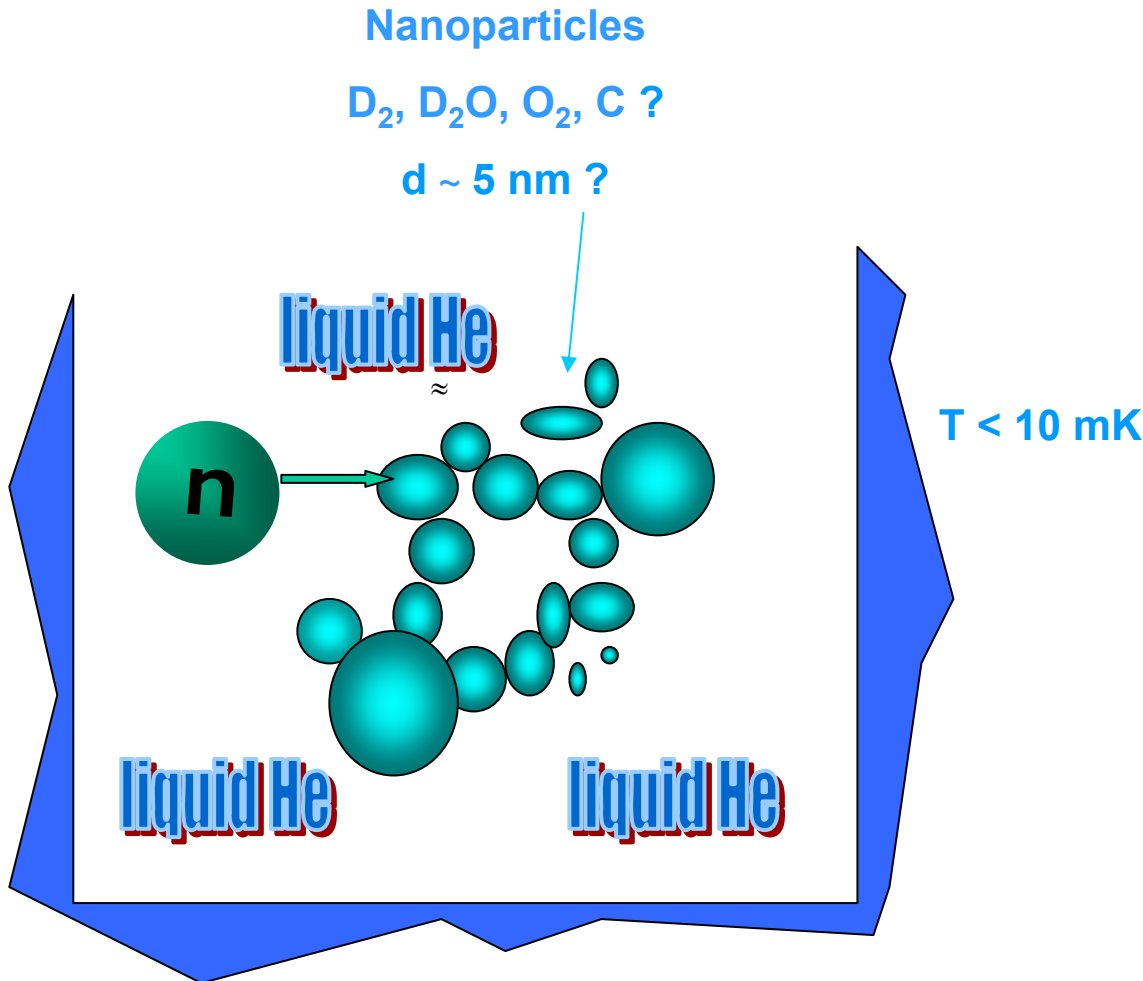
V.V.N., Physics of Atomic Nuclear 65(3), 400-408 (2002)

V.V.N., Interaction of neutrons with nanoparticles. *Physics of Atomic Nuclear*, 2002. 65(3): p. 400-408.



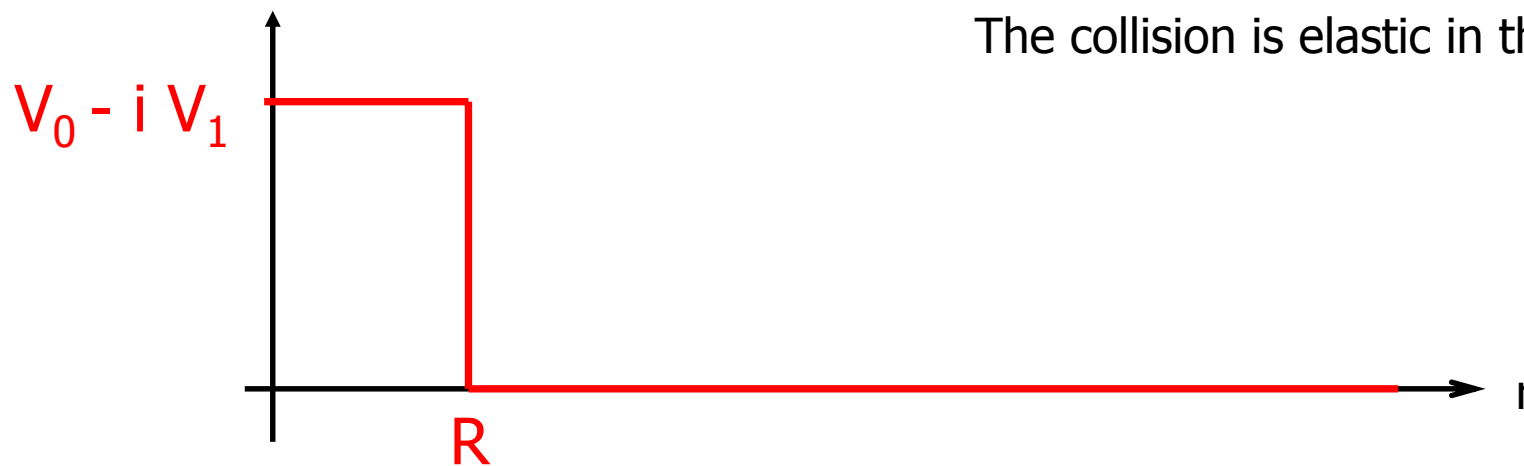
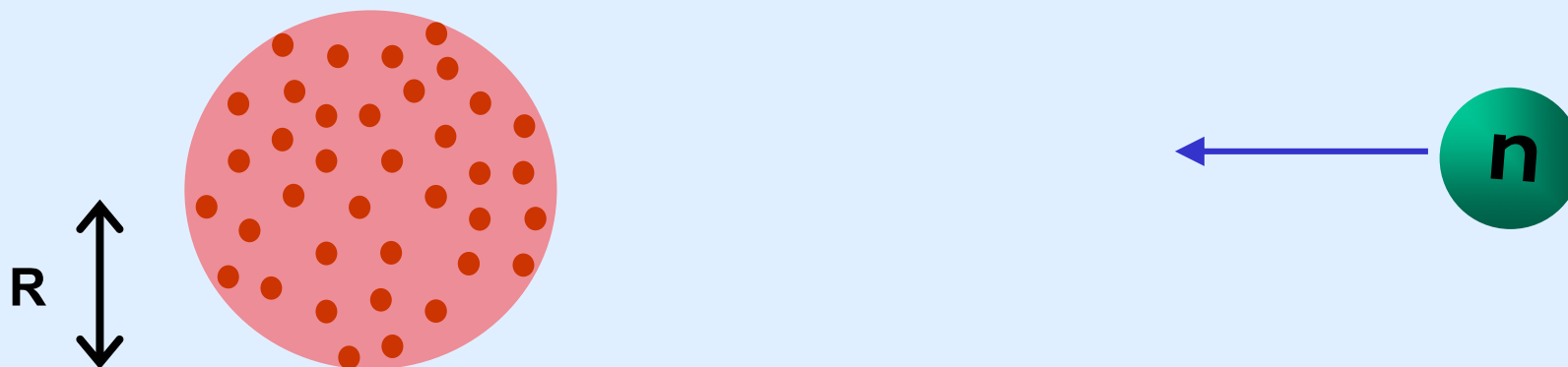
Motivations: 2) Reactor moderators

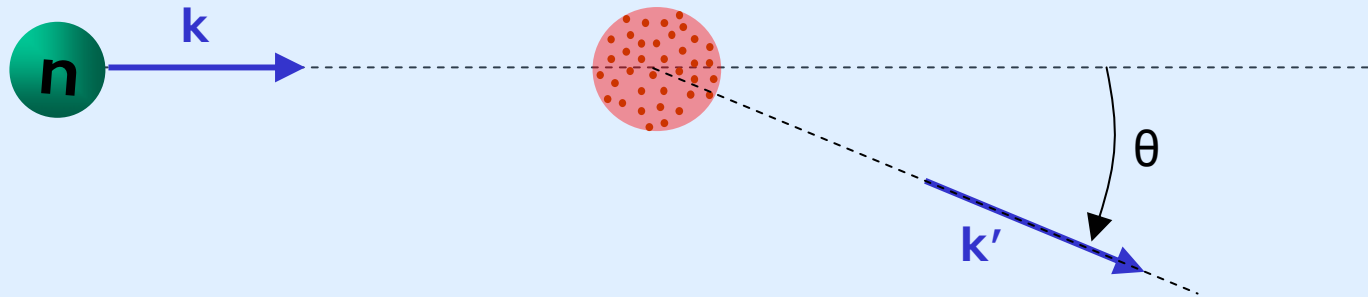




Neutrons can excite many degrees of freedom:

- Simple collision with a single nanoparticle
- Rotation
- Collective degrees of freedom (phonons, rotons)
- Breaking of inter-particle bonds





Three main quantities can be derived from the scattering amplitude $f(\theta)$:

- The diffusion cross-section

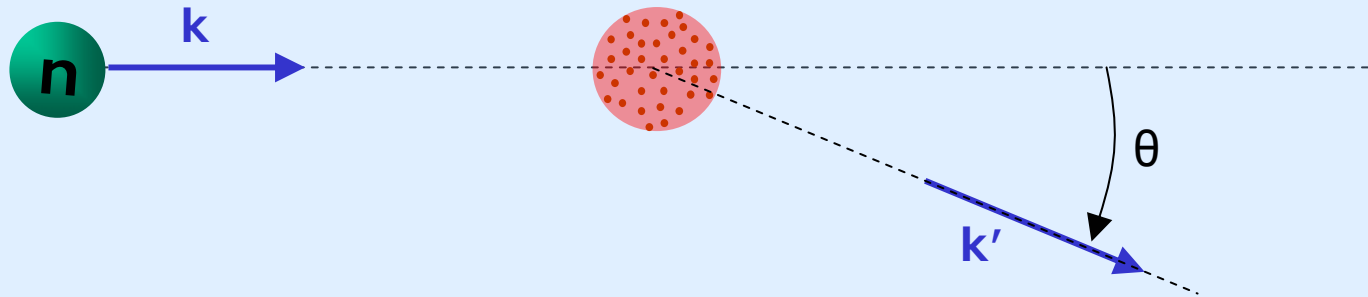
$$\sigma_s = \int |f|^2 d\Omega = 2\pi \int_0^\pi |f(\theta)|^2 \sin(\theta) d\theta$$

- The absorption cross-section

$$\text{Im}(f(\theta = 0)) = \frac{k}{4\pi} \sigma_a$$

- The relative mean energy loss per collision

$$\xi = \frac{\langle \Delta E \rangle}{E} = \frac{4}{A} \langle \sin^2(\theta/2) \rangle$$



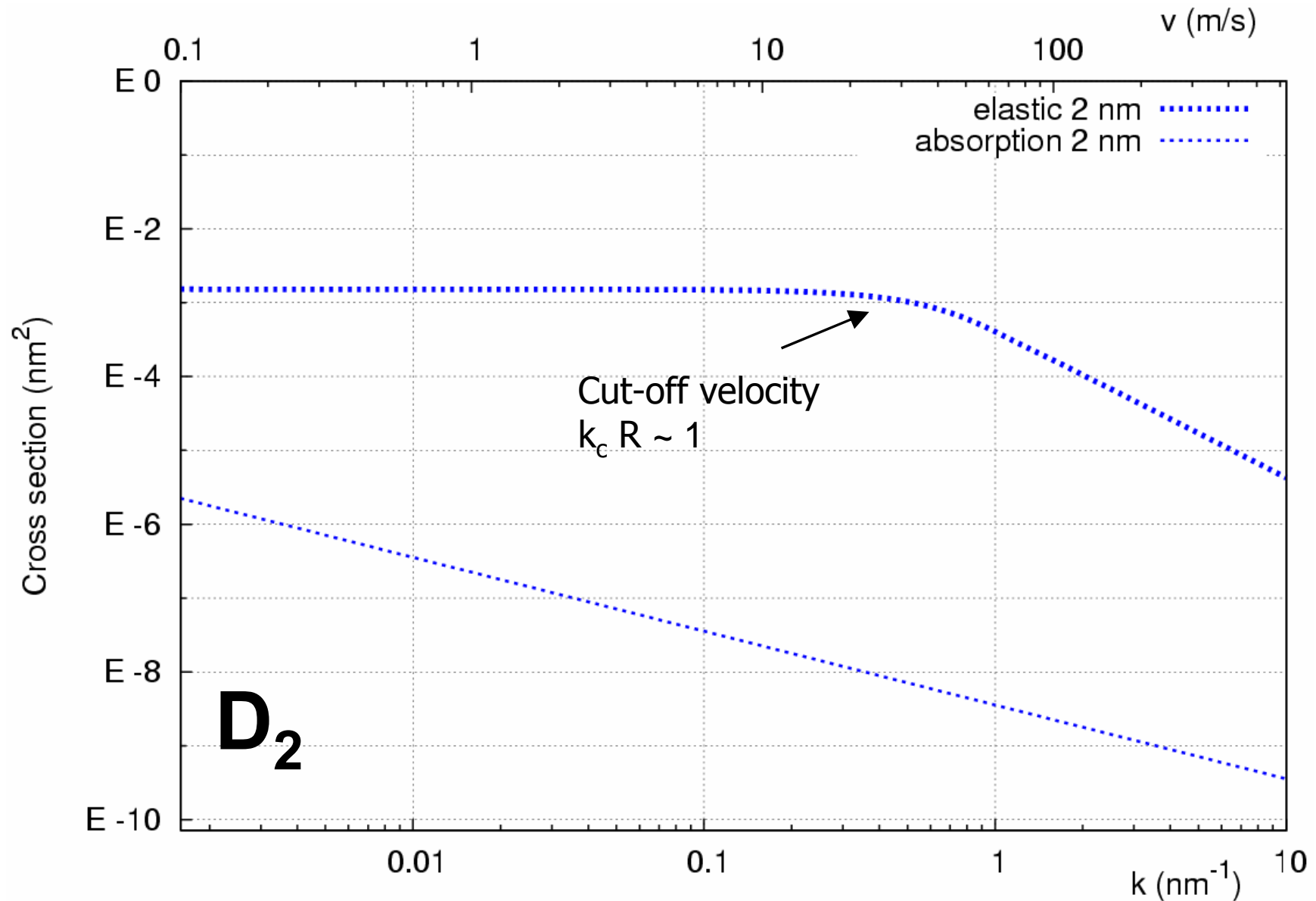
Calculations at the first order of the Born approximation

Validity : $VR^2 \ll \frac{\hbar^2}{2m}$

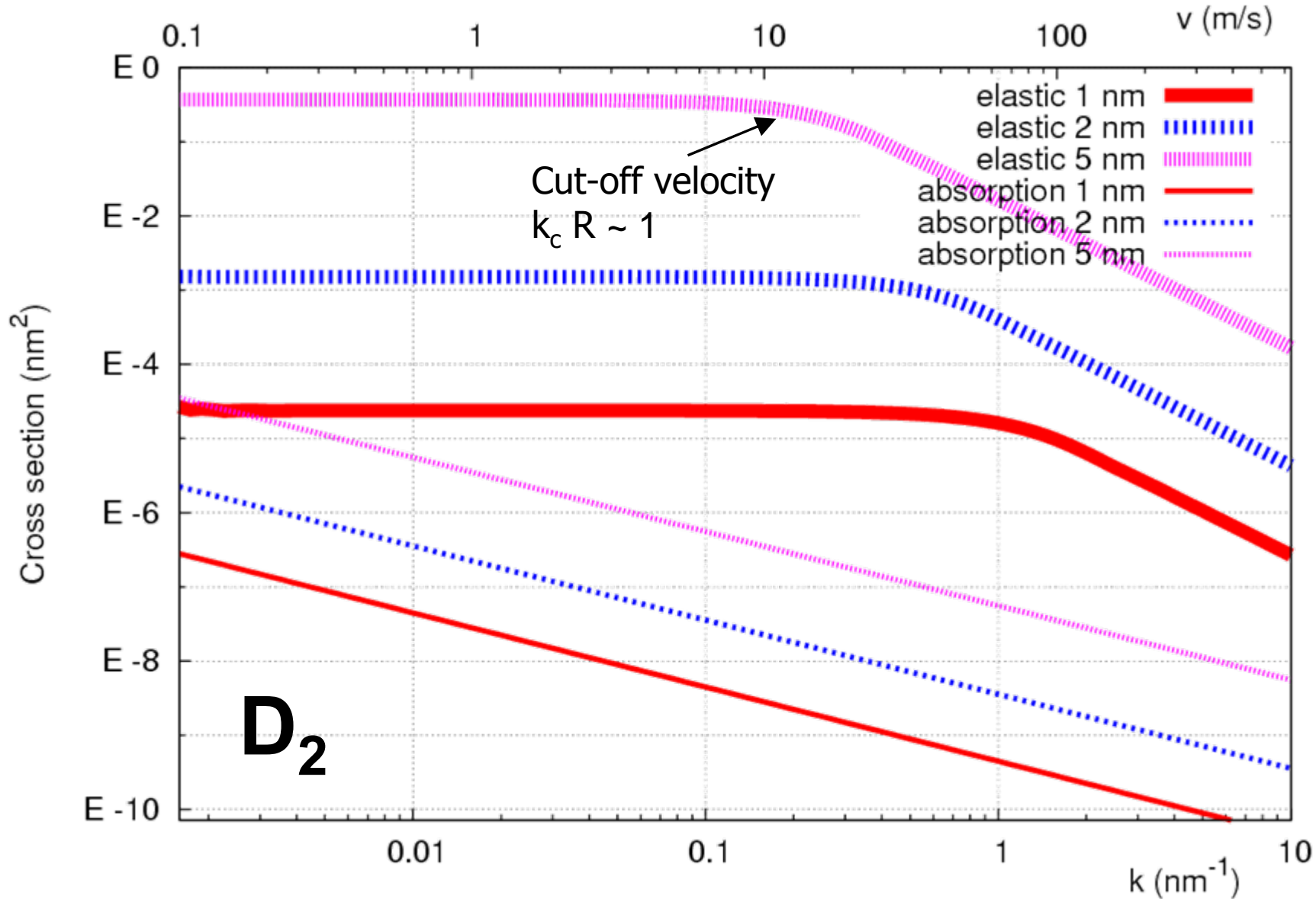
$$VR_{\text{Born}}^2 = \frac{1}{10} \frac{\hbar^2}{2m}$$

nanoparticle	D ₂	D ₂ O	O ₂	CO ₂	C (Diamond)	Be
R _{Born} (nm)	4.5	3.7	5.4	4.5	2.6	2.9

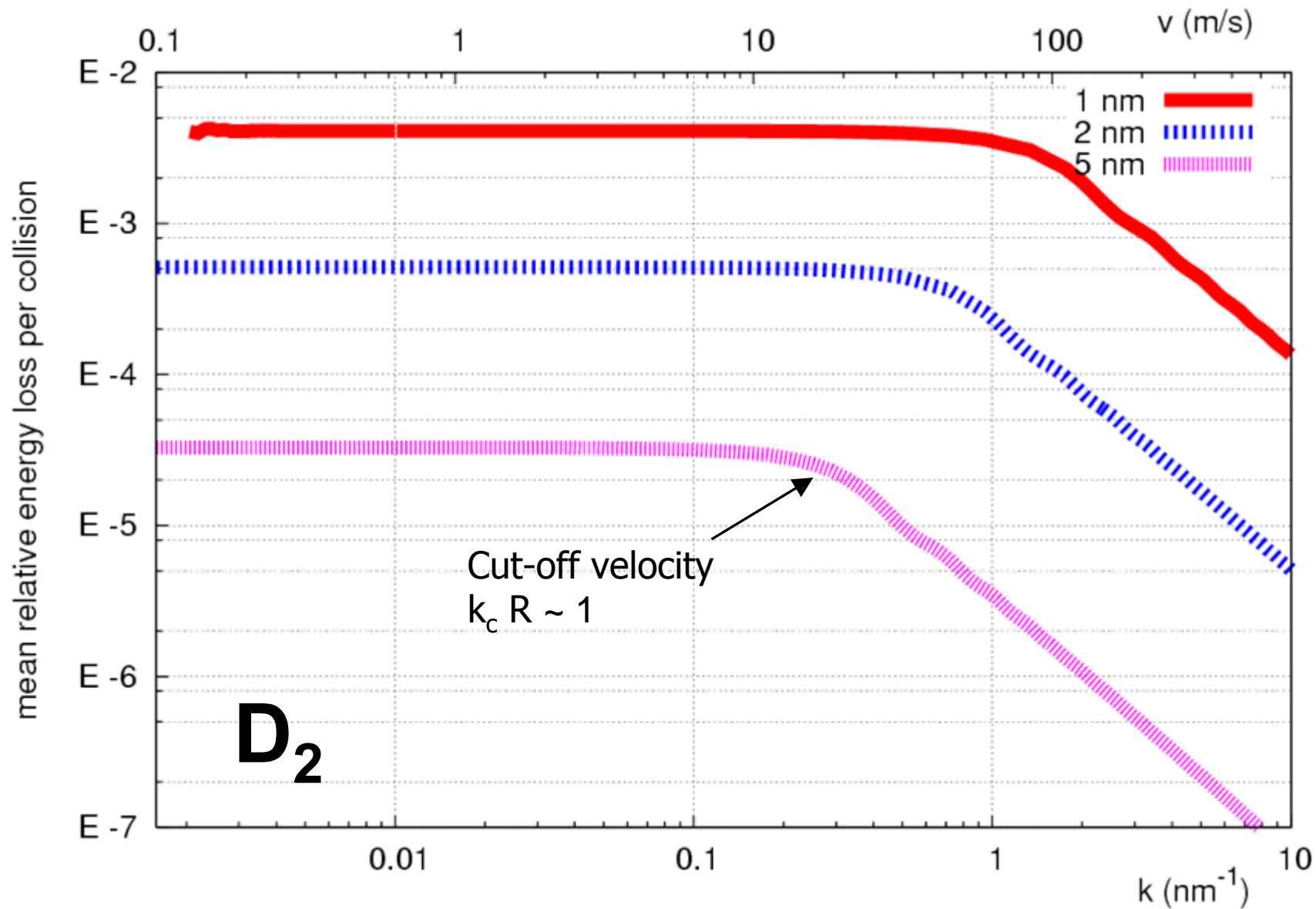
Result :
$$f(\theta) = -\frac{1}{4\pi} \frac{2m}{\hbar^2} \int e^{i(\mathbf{k}-\mathbf{k}')\cdot\mathbf{r}} V(\mathbf{r}) d\mathbf{r}$$



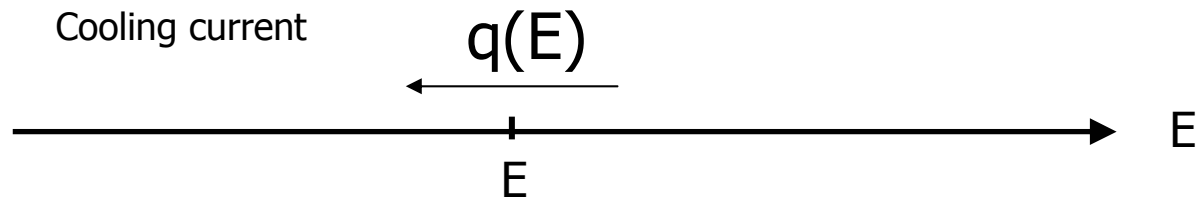
Dependence on radius



Energy loss



- Realistic model $n(\mathbf{r}, \mathbf{v}, t)$ solution of the (too complicated) Boltzmann equation
- Infinite Moderator $n(E, t)$



Equation of conservation for the number of neutrons

$$\frac{\partial n}{\partial t} - \frac{\partial q}{\partial E} + \text{Absorption} - \text{Source} = 0$$

Conservation equation

$$\frac{\partial n}{\partial t} - \frac{\partial q}{\partial E} + \text{Absorption} - \text{Source} = 0$$

Neutron flux

$$\phi(E) = v n(E) = \sqrt{\frac{2E}{m}} n(E)$$

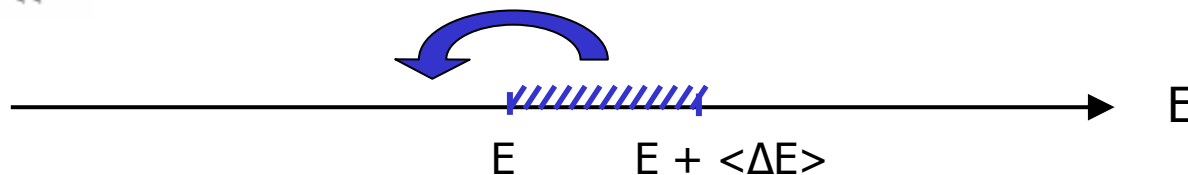
Absorption term

$$dN_a = \Sigma_a \phi(E) dE dt$$

Cooling current

$$q(E) = E \xi(E) \Sigma_s(E) \phi(E)$$

$$\xi \ll 1$$



Conservation equation

$$\frac{\partial n}{\partial t} - \frac{\partial q}{\partial E} + \text{Absorption} - \text{Source} = 0$$

With flux variable

$$\sqrt{\frac{m}{2E}} \frac{\partial \phi}{\partial t} - \frac{\partial q}{\partial E} + \Sigma_a \phi = \text{Source}$$

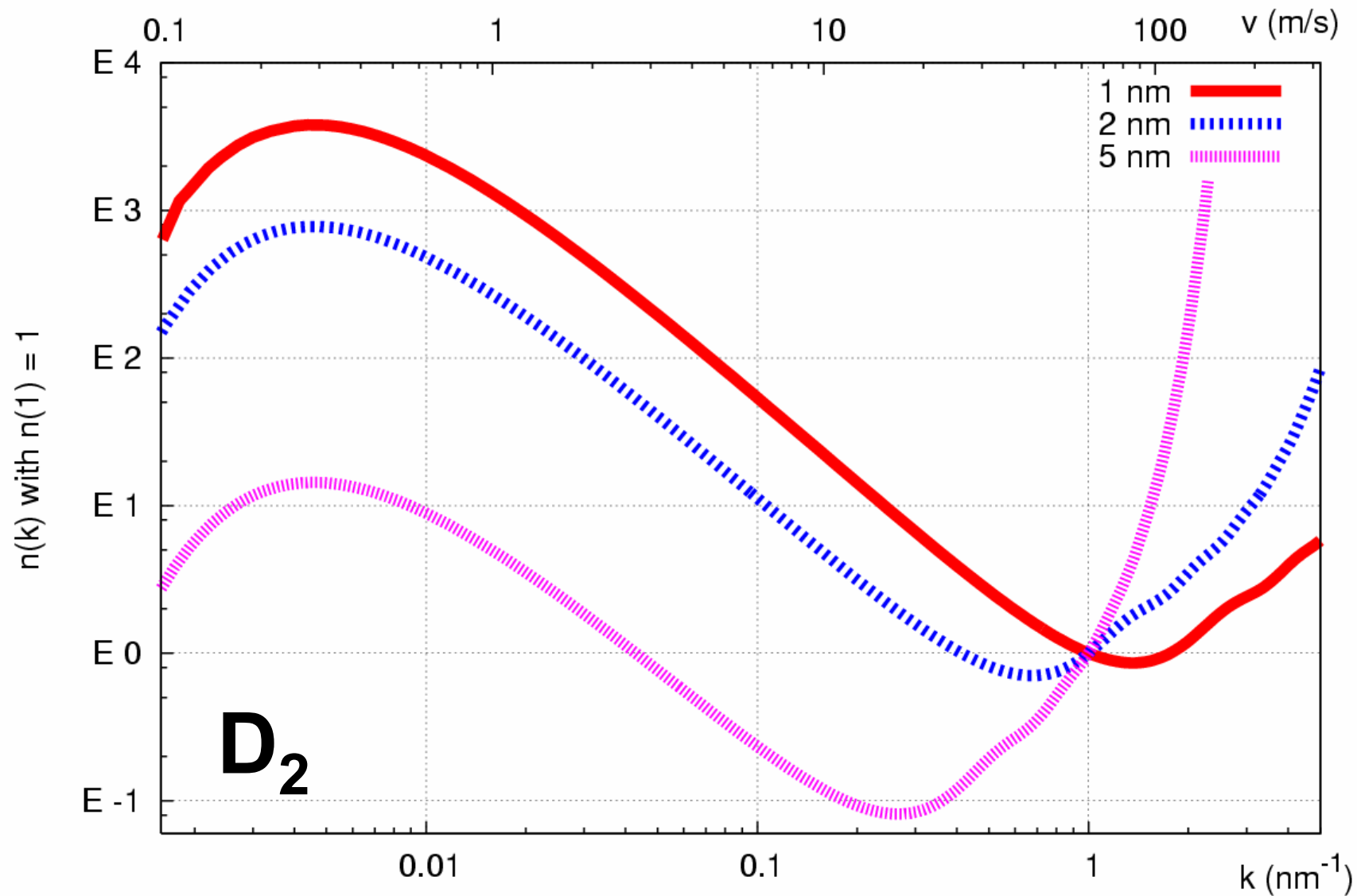
In stationary regime, with a mono-energetic source at energy E_0

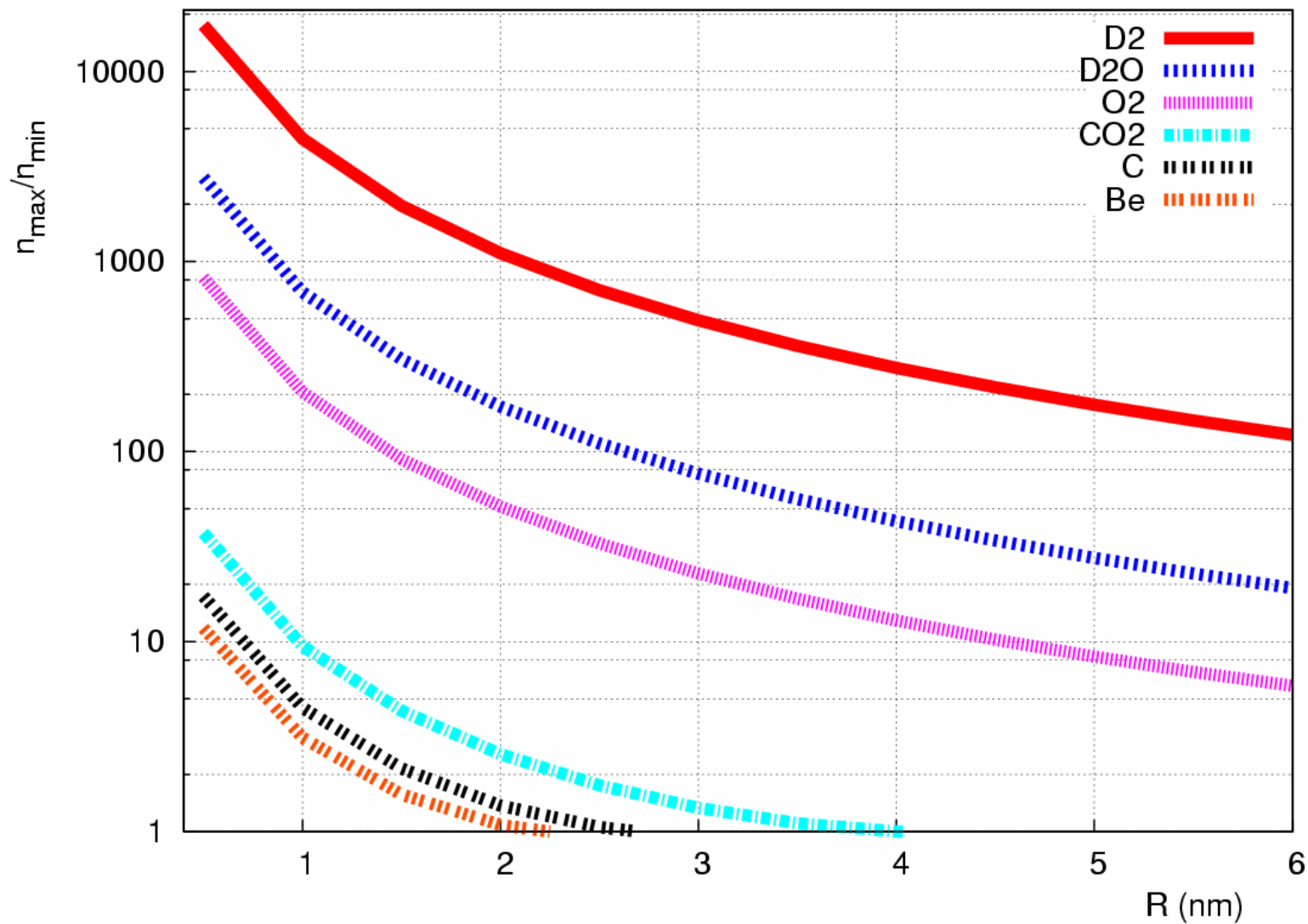
$$-\frac{d}{dE} (\xi \Sigma_s E \phi) + \Sigma_a \phi = 0$$

Solution:

$$n_{stat}(E) = n_{stat}(E_0) \left(\frac{E_0}{E} \right)^{3/2} \frac{\xi(E_0) \Sigma_s(E_0)}{\xi(E) \Sigma_s(E)} \exp \left(- \int_E^{E_0} \frac{\Sigma_a}{\xi \Sigma_s} \frac{d\epsilon}{\epsilon} \right)$$

Efficiency ranges

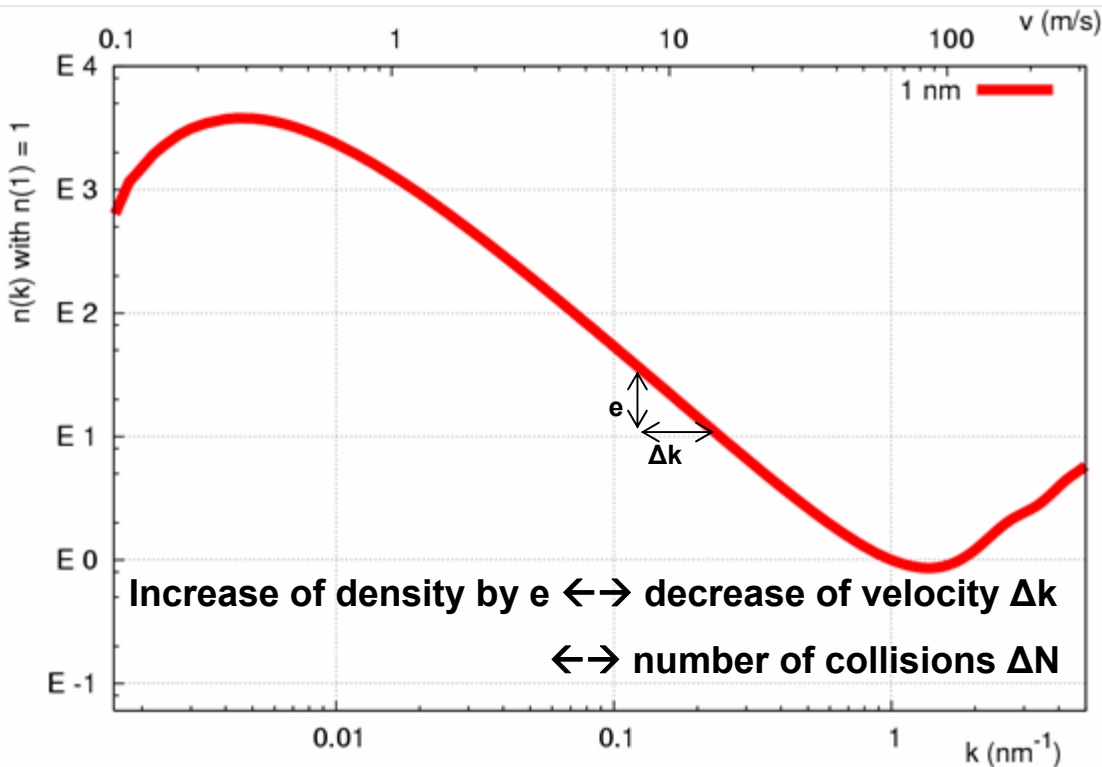




Maximum increase of density in the phase space for different materials.

Two important quantities to consider for a more realistic moderator

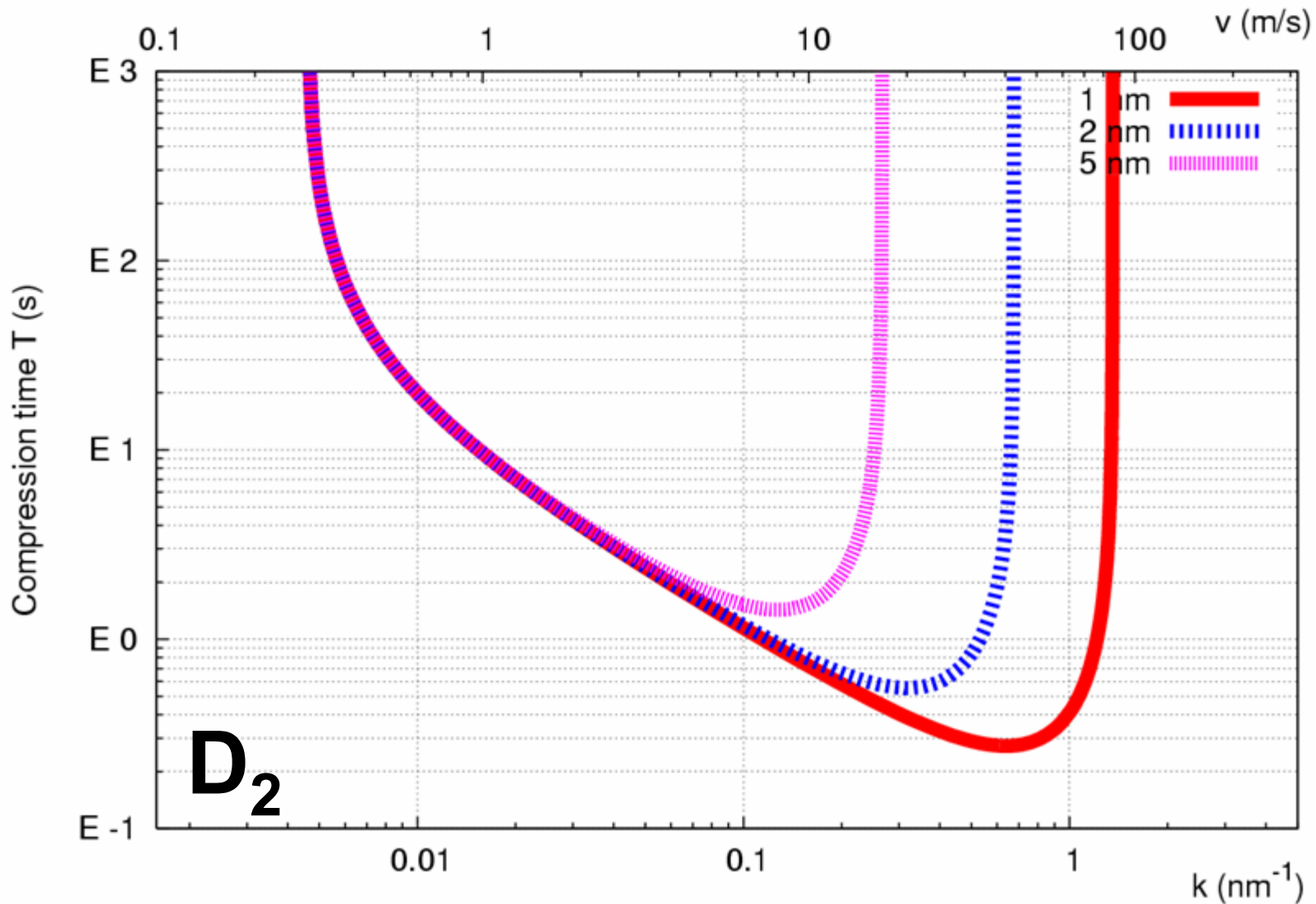
1. The Size of the moderator
2. The thermalization time



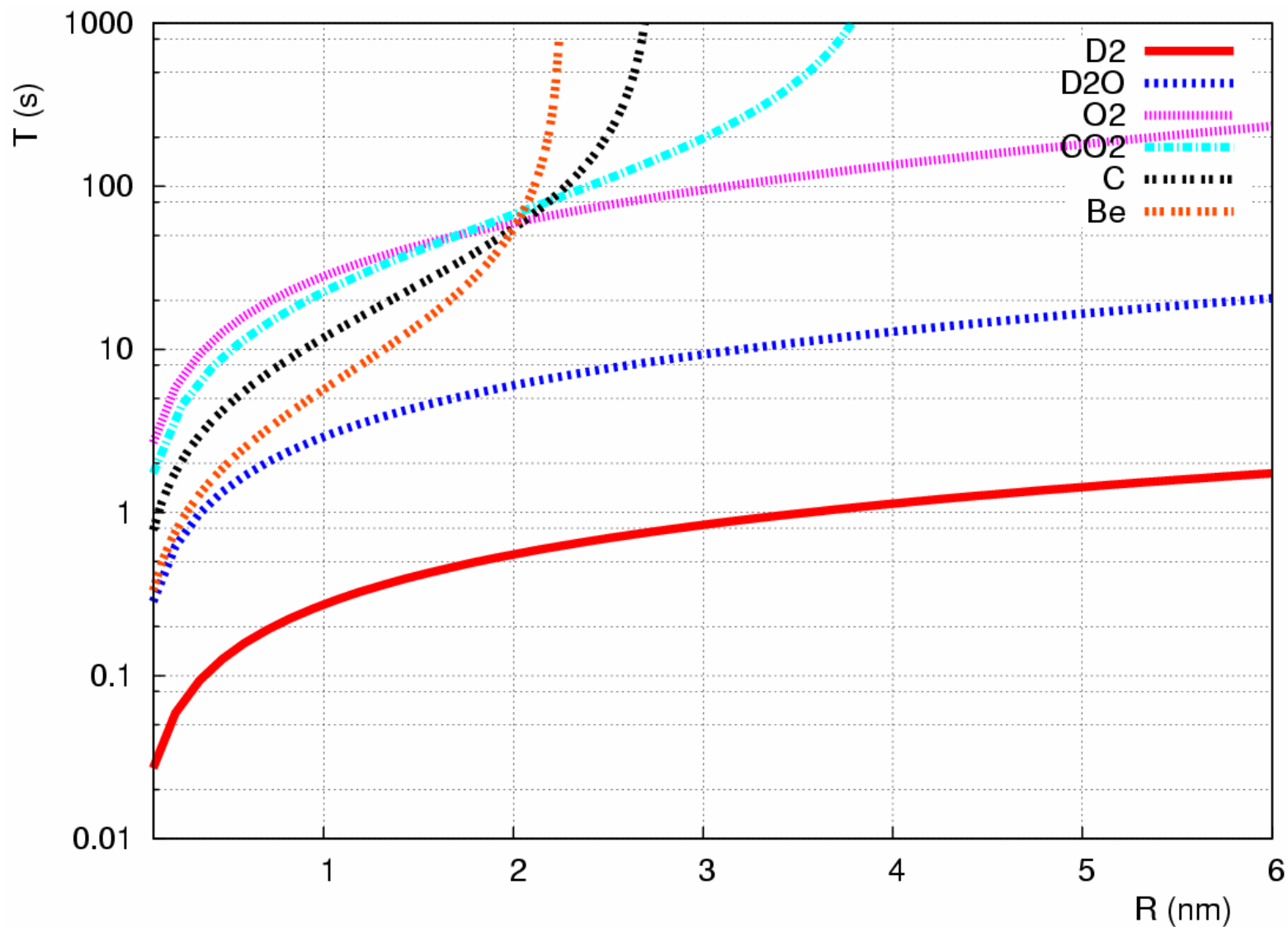
We can estimate, for ΔN collisions:

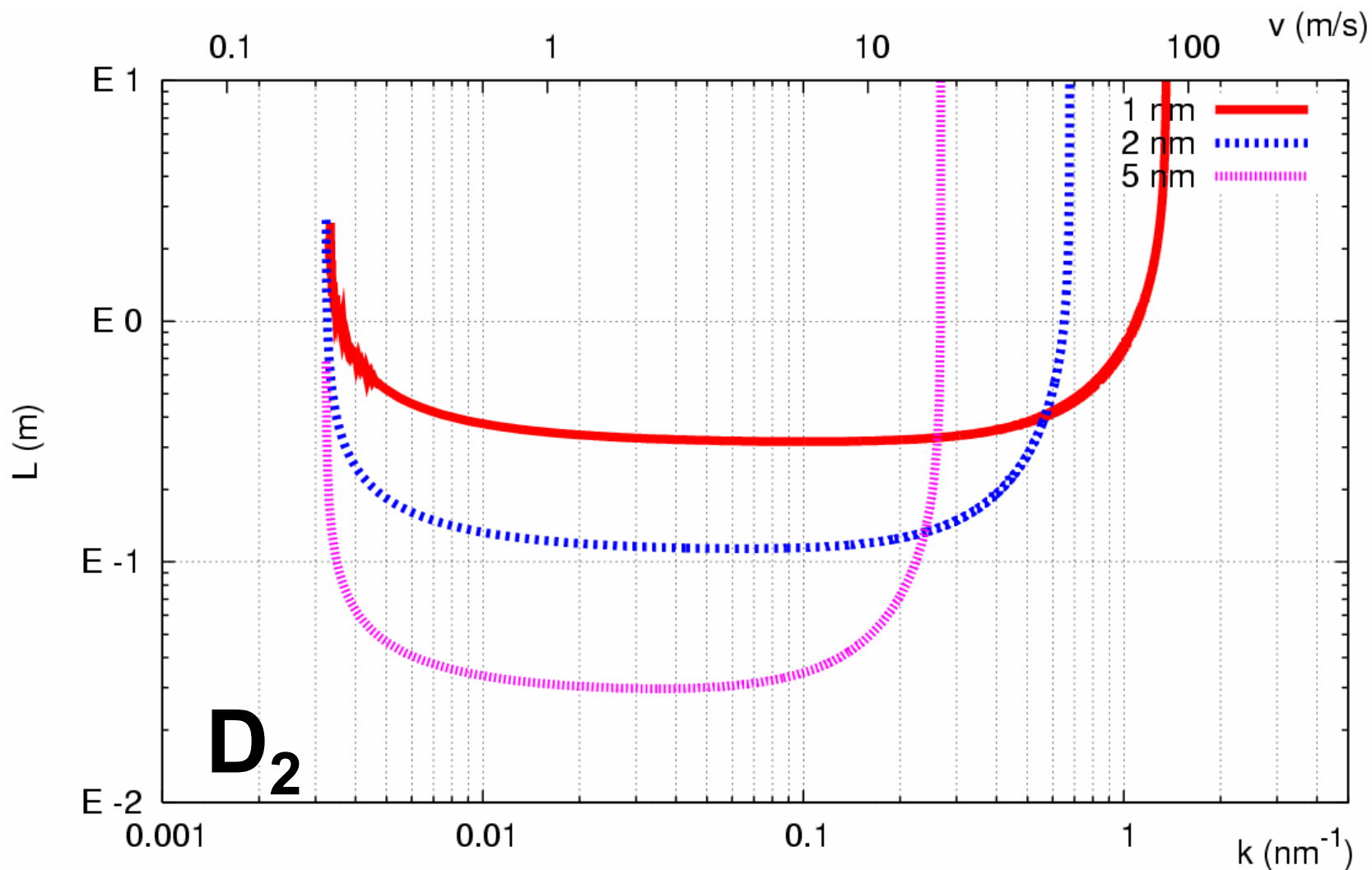
1. The square mean root of the distance traveled $L(k)$
2. The mean time it lasted $T(k)$

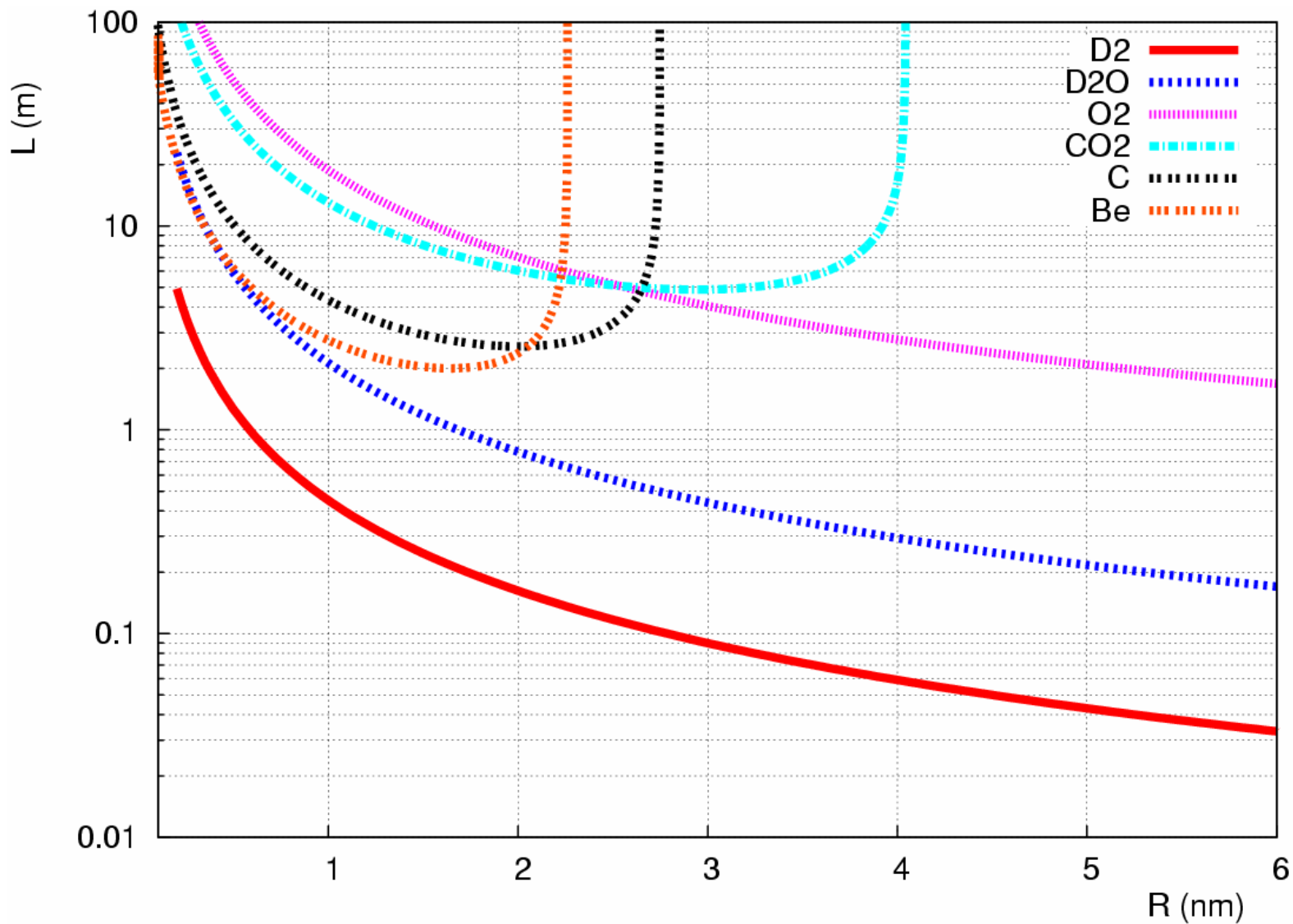
Compression time for different radius



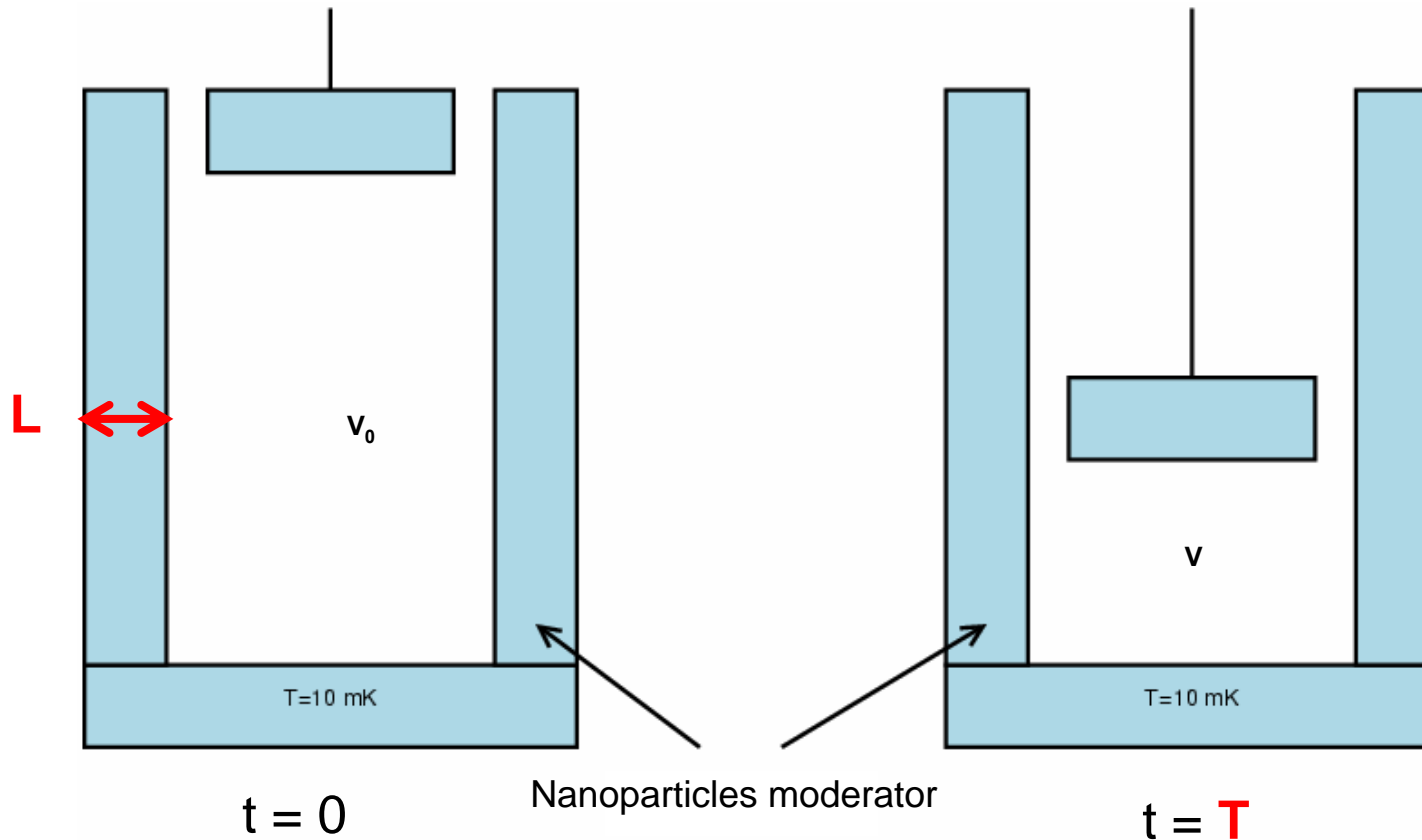
Compression time for different materials







Options for realistic moderators



There is no theoretical limitations on the increase of density

- **This model neglects any other degrees of freedom of nanoparticles in gels → further theoretical work is needed**
- **Use of the Born approximation → not a problem**
- **Limitation at high energy → not a problem**
- **Limitation at low energy → is there a problem ?**

- The best nanoparticle material is Deuterium.
- The range of efficient moderation is **~1-100 m/s**

Very cold neutrons

- The optimal size of the nanoparticles is about $R=2$ nm
- With the thermalisation time of a few seconds and the moderator size of a few times 10 cm, a principle realization of the presented cooling mechanism is feasible

